

MET3502/5561 Lecture 18: Potential Vorticity

Supplemental Reading: Holton's book, Section 4.3

Motivation: Need a simple equation to diagnose synoptic-dynamics that is easier to interpret than the vorticity equation and more complete than the barotropic vorticity equation.

- As shown in Holton's books section 4.3, assuming adiabatic & frictionless conditions, the potential vorticity (PV) is conserved following fluid motion:

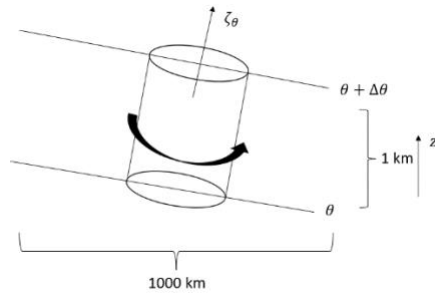
$$PV \equiv (\zeta_\theta + f) \left(-g \frac{\partial \theta}{\partial P} \right), \frac{DPV}{Dt} = 0 \text{ if adiabatic \& frictionless}$$

(potential temperature θ is conserved if adiabatic)

- Advantage: can be used for divergent (vertical motion $w \neq 0$) & baroclinic (density $\rho \neq \text{const}$) conditions, such as is found in the actual atmosphere.

Components of Potential Vorticity

- ζ_θ = relative vorticity in isentropic coordinates:
 - Isentropic coordinates use potential temperature (θ) to define horizontal surfaces.
 - Remember \rightarrow θ increases with height (unless super adiabatic)



- ζ_θ is simply the component of vorticity normal to the potential temperature surfaces.
 - Most of the time (But not always) $|\zeta_\theta| \approx |\zeta_z|$
- f = planetary vorticity
 - g = gravity [assume constant]
 - $-\frac{\partial \theta}{\partial P}$ = Static Stability

θ is potential temperature: $\theta = T \left(\frac{1000 \text{ mb}}{P} \right)^{R_d/C_p}$, where P is pressure, C_p is specific heat of dry air when pressure is constant, R_d is the individual gas constant for dry air.

$$\Rightarrow \frac{d\theta}{dz} = \left(\frac{1000}{P} \right)^{R_d/C_p} \frac{dT}{dz} + T \cdot 1000^{R_d/C_p} \frac{dP}{dz} \left(-\frac{R_d}{C_p} \right)$$

Since $\frac{\theta}{T} = \left(\frac{1000}{P} \right)^{R_d/C_p}$ So:

$$\begin{aligned} \frac{d\theta}{dz} &= \frac{\theta}{T} \frac{dT}{dz} - \frac{R_d}{C_p} T \frac{1000^{R_d/C_p}}{P^{(R_d/C_p+1)}} \frac{dP}{dz} \\ &= \frac{\theta}{T} \frac{dT}{dz} - \frac{R_d}{C_p} \theta \cdot \frac{1}{P} \frac{dP}{dz} \end{aligned}$$

From Hydrostatic Equation: $dP = -\rho g dz$, $\rho = \frac{P}{R_d T}$

$$\Rightarrow \frac{dP}{dz} = -\rho g = -\frac{Pg}{R_d T}$$

$$\text{So } \frac{d\theta}{dz} = \frac{\theta}{T} \frac{dT}{dz} + \frac{R_d \theta}{C_p P} \frac{Pg}{R_d T} = \frac{\theta}{T} \frac{dT}{dz} + \frac{\theta}{T} \frac{g}{C_p}$$

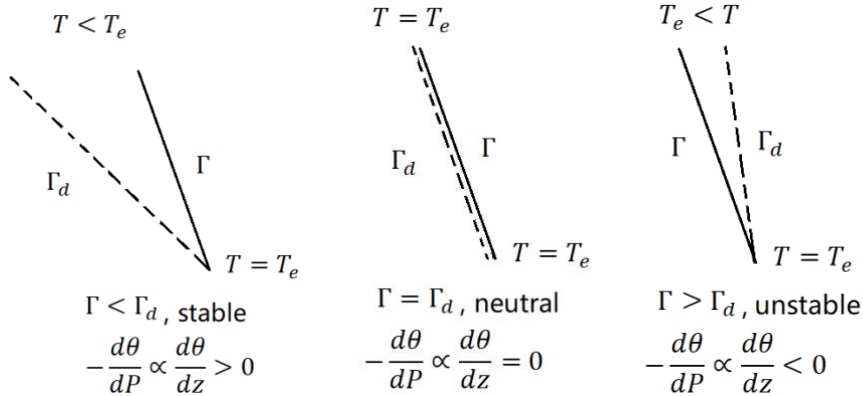
$$\frac{d\theta}{dz} = \frac{\theta}{T} \left(\frac{dT}{dz} + \frac{g}{C_p} \right)$$

Since: Dry adiabatic lapse rate $\Gamma_d = \frac{g}{C_p}$

Environment lapse rate $\Gamma = -\frac{dT}{dz}$

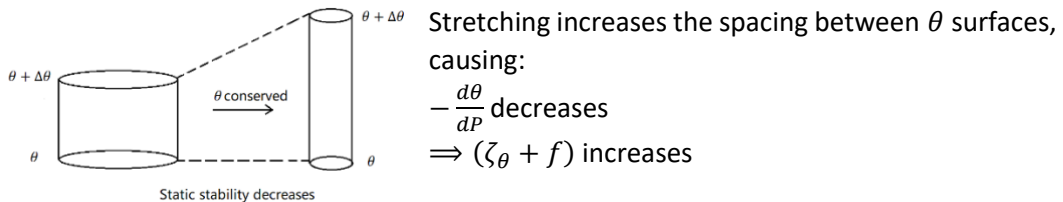
So: $\frac{d\theta}{dz} = \frac{\theta}{T} (\Gamma_d - \Gamma)$ since $-\frac{d\theta}{dP} \propto \frac{d\theta}{dz}$

$\Rightarrow -\frac{d\theta}{dP} \propto (\Gamma_d - \Gamma)$, which is the stability for dry air:



Physical Interpretation:

- PV is a function of:
 - 1) Absolute Vorticity $\zeta_\theta + f$
 - 2) Static Stability $\left(-\frac{d\theta}{dP}\right)$
- If PV is conserved, then an increase of either $(\zeta_\theta + f)$ or $-\frac{d\theta}{dP}$ must lead to a decrease of the other (or vice versa).
- If a fluid column is stretched, the absolute vorticity must increase because:



- If a fluid column is squished, the absolute vorticity must decrease because:

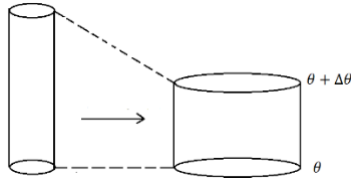
Squishing decreases the spacing between θ surfaces, causing: $-\frac{d\theta}{dP}$ increases $\Rightarrow (\zeta_\theta + f)$ decreases



Applications of PV:

- Example 1: What happens when an air column moves from a region of low static stability to a region of high static stability?

- If static stability is increasing, then $-\frac{d\theta}{dP}$ is increasing \Rightarrow the θ surface must be getting closer together:

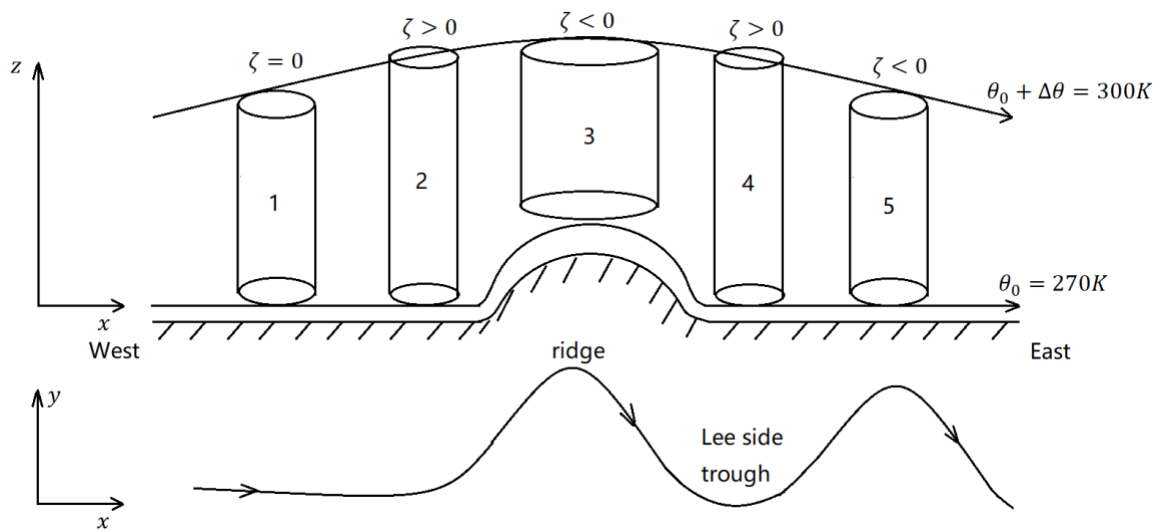


Column is squished \rightarrow parcel must decrease the absolute vorticity.

Since $PV = g(\zeta_\theta + f) \left(-\frac{d\theta}{dP}\right) = \text{constant}$,

$-\frac{d\theta}{dP}$ increases $\Rightarrow (\zeta_\theta + f)$ must decrease.

- Example 2 (Holton's book, Page 100-102): Zonal westerly flow with no initial relative vorticity impinging on a mountain barrier — Lee cyclone genesis



Position	$-\frac{d\theta}{dP}$	f	$\zeta_\theta (\sim \zeta)$
1	> 0	f_0 (doesn't change much)	$= 0$
1-2	Decreasing (stretching)	f_0	must increase
2		f_0	$\zeta > 0$
2-3	Increasing (squishing)	f_0	must decrease
3		f_0	$\zeta < 0$
3-4	Decreasing (stretching)	f_0	must increase
4		f_0	$\zeta > 0$
4-5	Increasing (squishing)	f_0	must decrease
5		f_0	$\zeta < 0$

Result: Steady westerly flow over a large-scale mountain barrier will result in a cyclonic flow pattern immediately to the east of the barrier (the lee side trough) followed by an alternating series of ridges & troughs.

Please read Holton's book, Page 99-102 for easterly flow example.

- Example 3: What is the impact of equatorward motion of an air column that has constant relative vorticity.

$$PV = -g \left(\frac{d\theta}{dP} \right) (\zeta_\theta + f) = \text{constant}$$

$g = \text{constant} \ \& \ \zeta_\theta = \text{constant} \ (\text{for this example})$

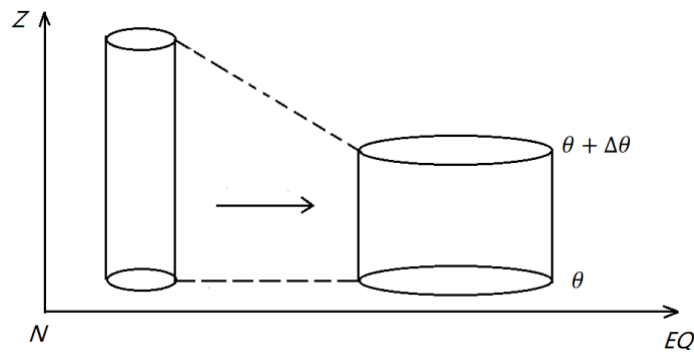
$$\begin{array}{c} \xrightarrow{A} \quad N \\ \xrightarrow{B} \quad EQ \end{array} \Rightarrow \left(-\frac{d\theta}{dP} \right)_A f_A = \left(-\frac{d\theta}{dP} \right)_B f_B = \text{constant}$$

As parcel moves equatorward, f decreases.

$\Rightarrow f_B < f_A$, and $-\frac{d\theta}{dP}$ must increase. So that

$$\left(-\frac{d\theta}{dP} \right)_B > \left(-\frac{d\theta}{dP} \right)_A \rightarrow \text{stability increases}$$

$\rightarrow \theta$ contours move closer together.



\Rightarrow If no change in relative vorticity is allowed, cold air moving equatorward will sink & warm.

Synoptic Application: Air moving equatorward will sink & warm if no relative vorticity is generated, while air moving poleward will rise & cool.