

Light Front (LF) Nuclear Structure

Gerald A Miller, U. of Washington

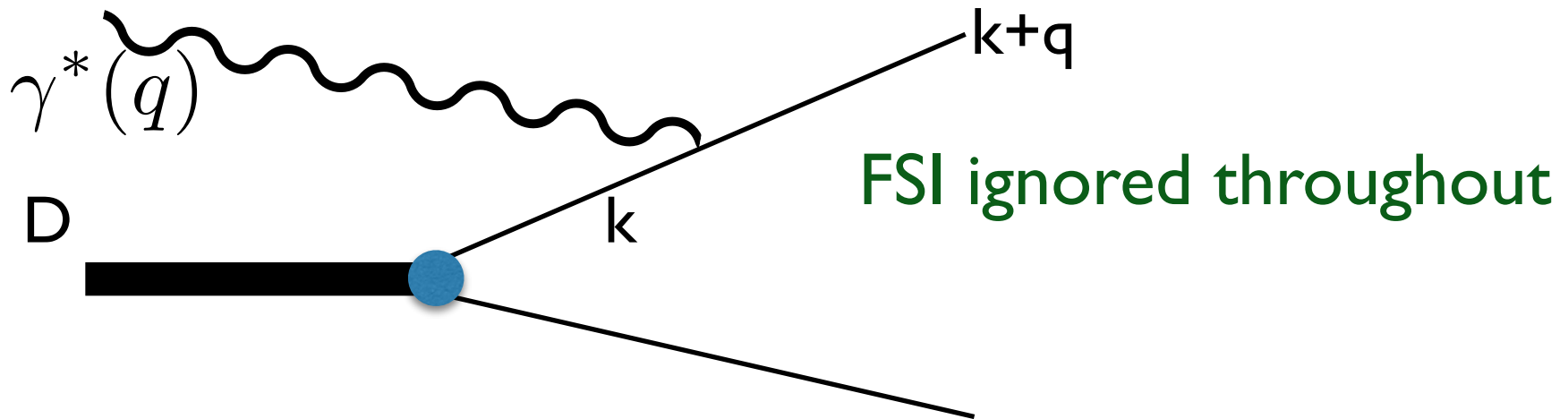
- formulate the NN interaction on the light front
- solve the Weinberg equation for the deuteron
- prescription from FS 81 review constructs LF wave function from NR wf:
- how good is this approx at recoil momenta few hundred MeV?
- can we get the LF wf from NN potentials?
- Nuclear Many Body LF wave function inc. Correlations -Miller & Machleidt
PRC60, 035202
- Relevance of LF wave functions for quasi-elastic scattering
stimulated by Ellie Long Proposal on Tensor asymmetries
work with John Terry REU
J R Cooke nucl-th/0112029 , Cooke & Miller PRC66, 034002
Miller, Prog. Nuc. Part. Phys. 45, 83
Tiburzi & Miller PRC81,035201

Light front quantization, Infinite momentum frame

“Time”, $x^+ = x^0 + x^3$, “Evolve”, $p^- = p^0 - p^3$

“Space”, $x^- = x^0 - x^3$, “Momentum”, p^+ (Bjorken)

Transverse position, momentum \mathbf{b}, \mathbf{p} $p^- = \frac{\mathbf{p}^2 + m^2}{p^+}$



If Photon energy $\gg m$ and struck nucleon \approx on mass – shell :

$$(k + q)^2 = m^2 \rightarrow k^+ q^- \approx Q^2, \quad \frac{Q^2}{\nu^2} \ll 1$$

FS '81

Integrate over k^-

$$d\sigma \sim \Psi_D^2(\mathbf{k}, \frac{k^+}{P_D^+} \equiv \alpha)$$

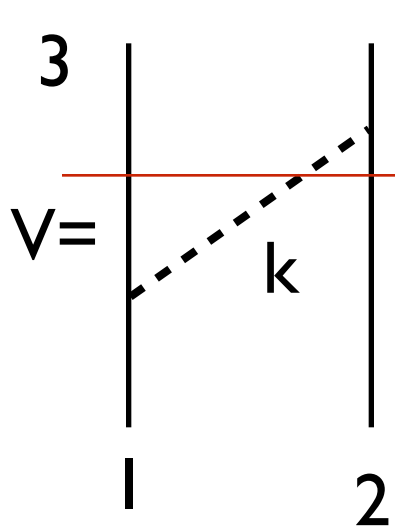
Are approximations valid at JLab 12- NO

Light front quantization, Infinite momentum frame

P^- is LF Hamiltonian, get from Lagrangian.

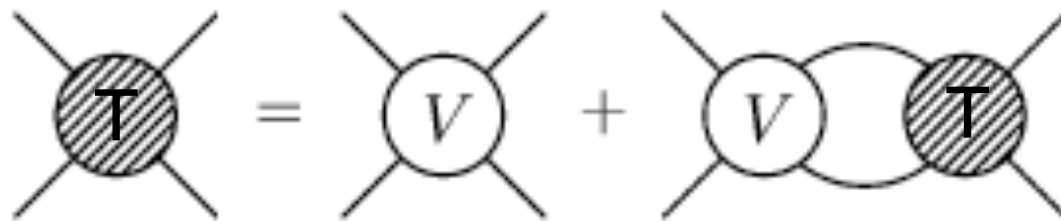
LF Schroedinger eq. $P^- |\Psi_D\rangle = M_D |\Psi_D\rangle$ Rest frame

One boson exchange



4 Usual Feynman diagram with recoil

Solve

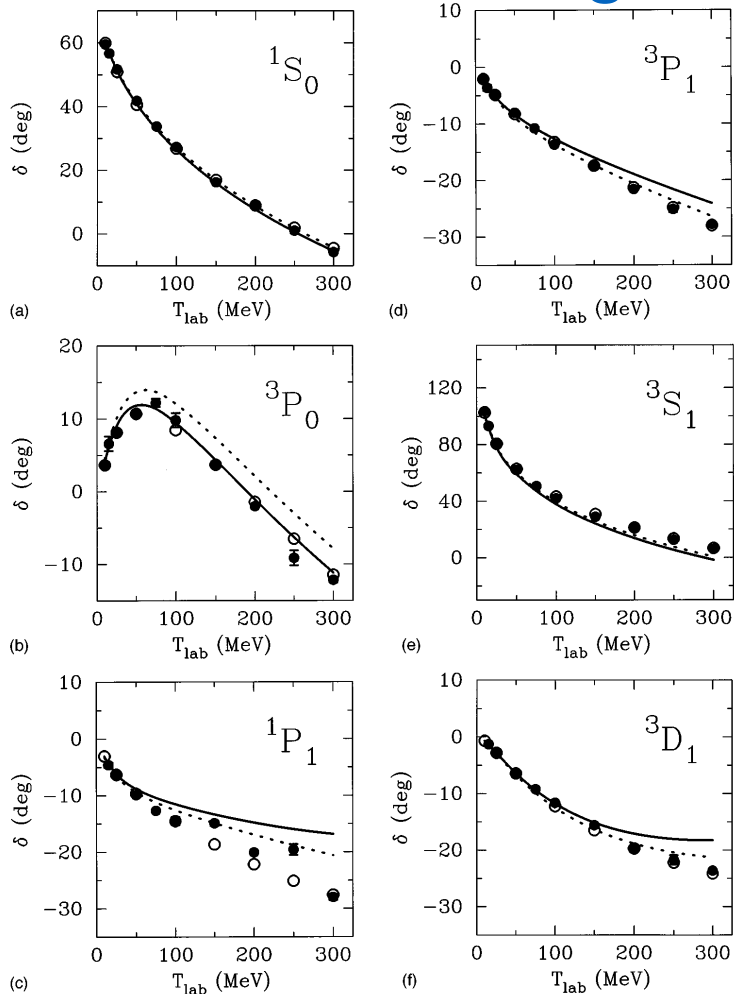


Weinberg equation = Lippmann-Schwinger eq with extra factor ($\sim I$) in Green's function

Miller & Machleidt PRC 60,035202

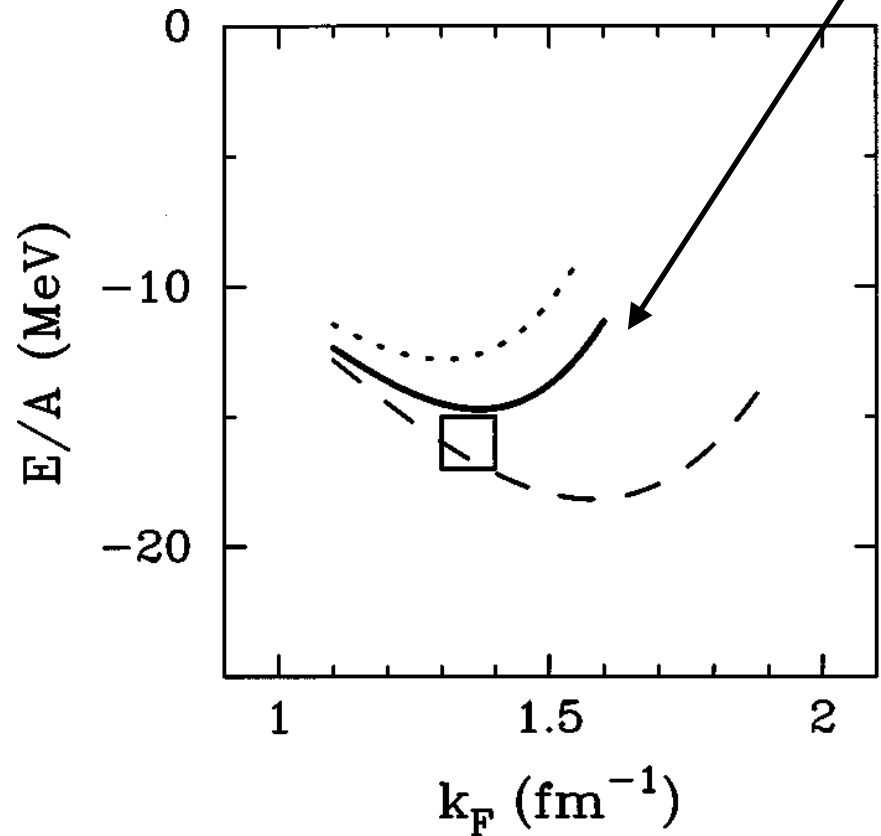
$\pi, \eta, \rho, \omega, a_0, \sigma$ exchange with extra factor in G

NN scattering



$$B_D = 2.245 \text{ MeV}$$

Nuclear Matter Saturation



The real problem- Bethe Salpeter Eq. (BSE)

$$T = \text{diagram} = \text{diagram } K + \text{diagram } K \text{ and } G$$

The diagram shows the Bethe-Salpeter equation: $T = K + K G$. The first term is a shaded circle with four external lines. The second term is a circle labeled 'K' with four external lines. The third term is a circle labeled 'K' connected to another shaded circle with four external lines, with a circle labeled 'G' between them.

K is sum of irred. diagrams

G depends on 4-momenta-product of two Feynman propagators

Reduce to 3 dimensions:

ET: integrate over k^0 . Ignore k^0 except in G . Sets relative time to 0.

LF: Integrate over k^- . Ignore k^- except in G . Sets relative $\tau = 0$

3 dimensional version of G is g_{ET} (Blankenbecler Sugar) or g_{LF} (Weinberg)

Puts BOTH particles on mass shell

No relation between wave functions in principle

Spectator on-shell- Gross equation- one particle off mass shell $k^2 - m^2 \neq 0$

Relation between equal-time and light-front wave functions

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The relation between equal-time and light-front wave functions is studied using models for which the four-dimensional solution of the Bethe-Salpeter wave function can be obtained. The popular prescription of defining the longitudinal momentum fraction using the instant-form free kinetic energy and third component of momentum is found to be incorrect except in the nonrelativistic limit. One may obtain light-front wave functions from rest-frame, instant-form wave functions by boosting the latter wave functions to the infinite momentum frame.

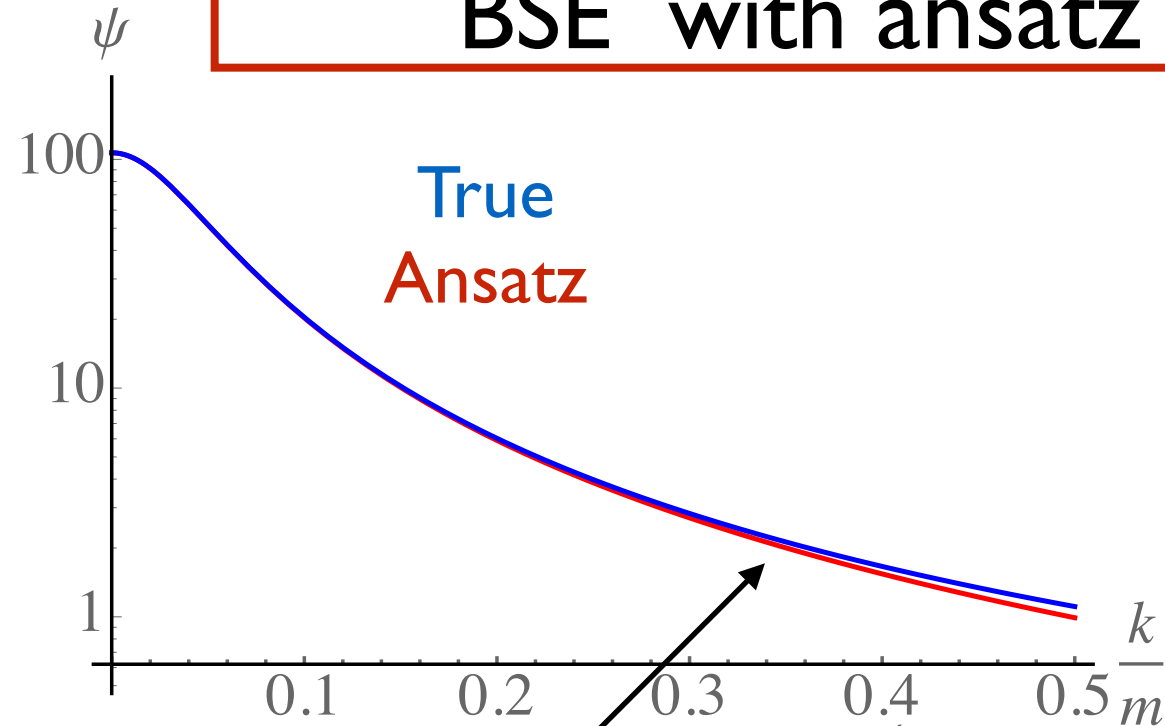
- How bad is the problem?
- Is D non-relativistic?
- Is ${}^3\text{He}$ non-relativistic?
- Answer by using solutions of Bethe-S eqn.



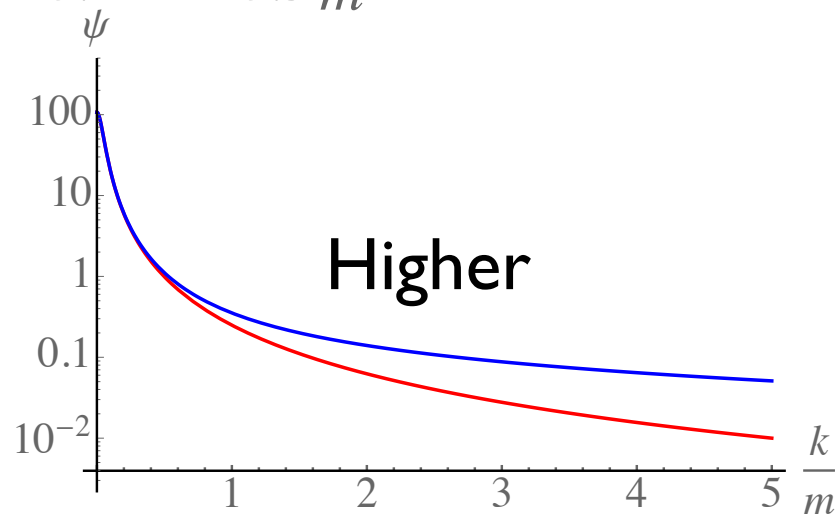
Model

FIG. 2. Bethe-Salpeter equation for a point interaction. The state is bound by the infinite chain of bubbles.

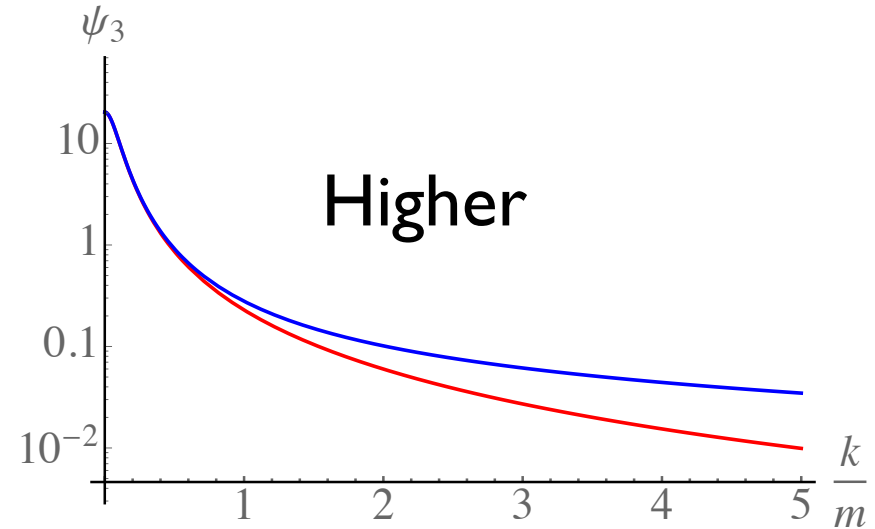
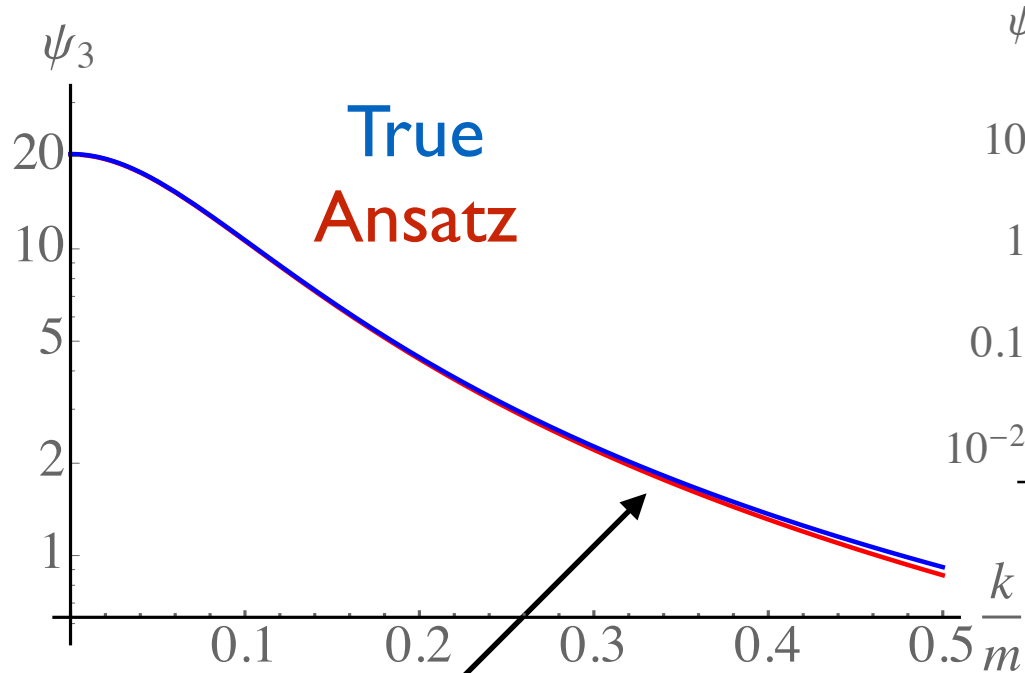
Deuteron Compares true LFD from BSE with ansatz from ET



OK up to here



^3He Compare true LFD from BSE with ansatz from ET

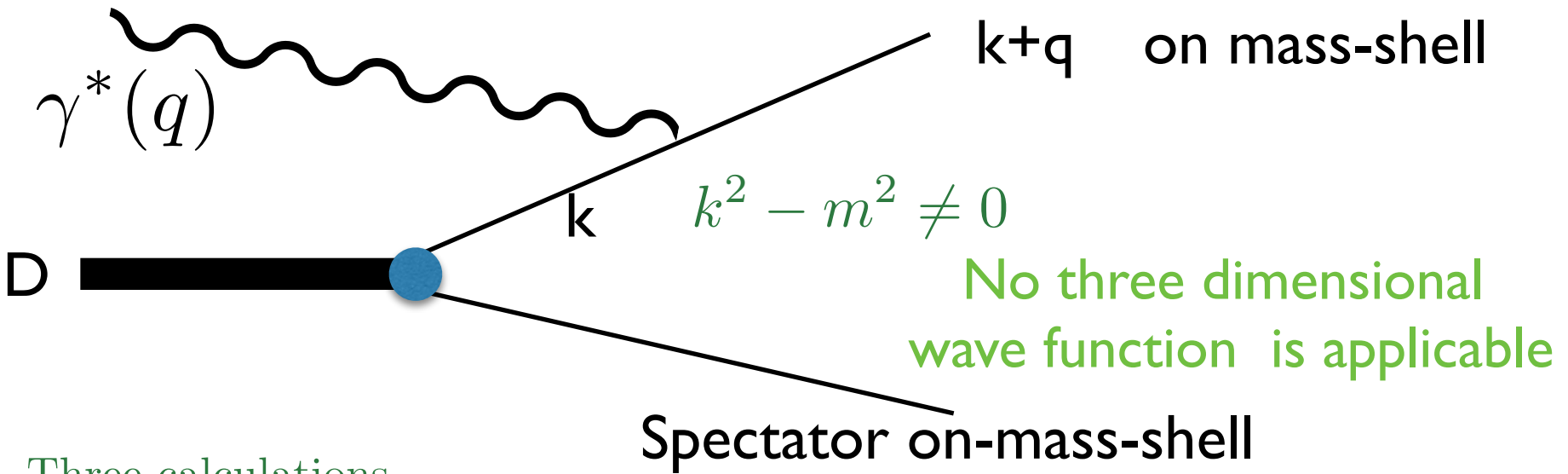


OK up to here

Are light front or equal time wave functions relevant?

- Light front wave functions - on-mass shell $k_i^2 = m^2$
- Same for equal time wave functions **Mismatch!**
- Answer this question through explicit example - point interaction model used above
- Space-like Form factors -use of light front wave functions gives exact results- GAM Phys.Rev. C80 (2009) 045210
- **What about quasi-elastic scattering?**

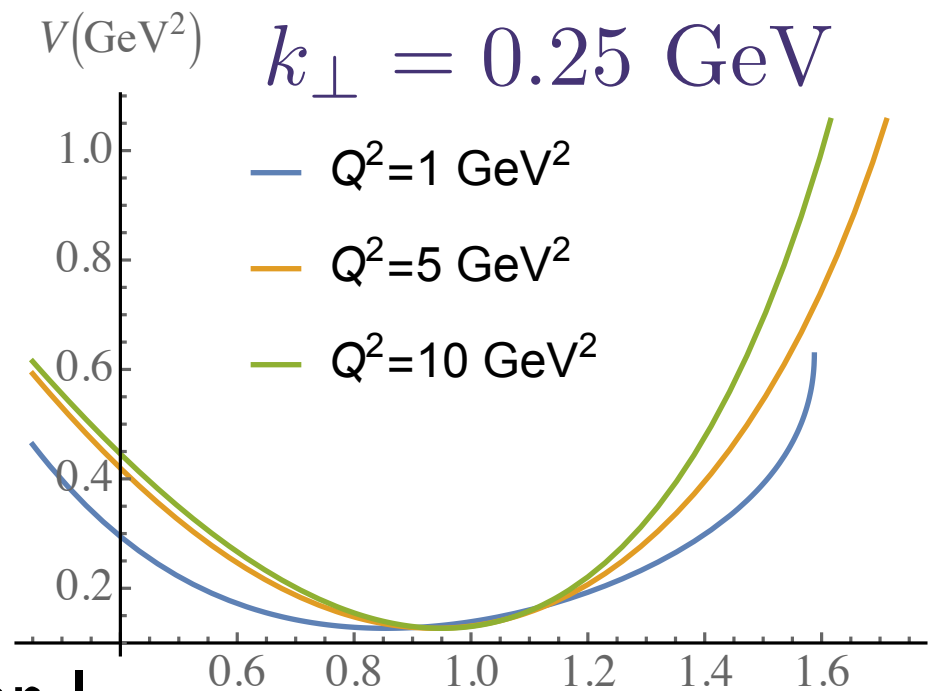
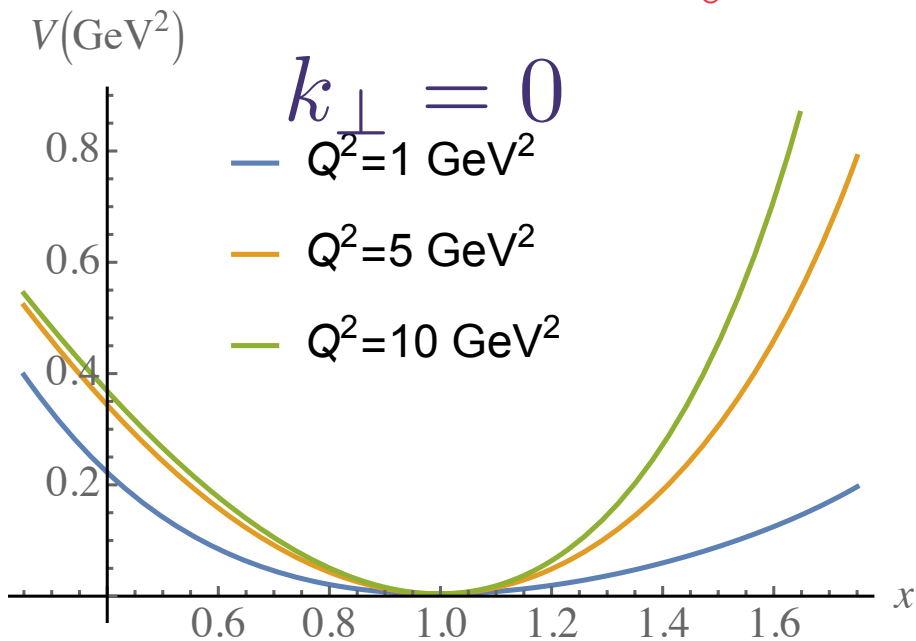
quasi-elastic scattering



Three calculations

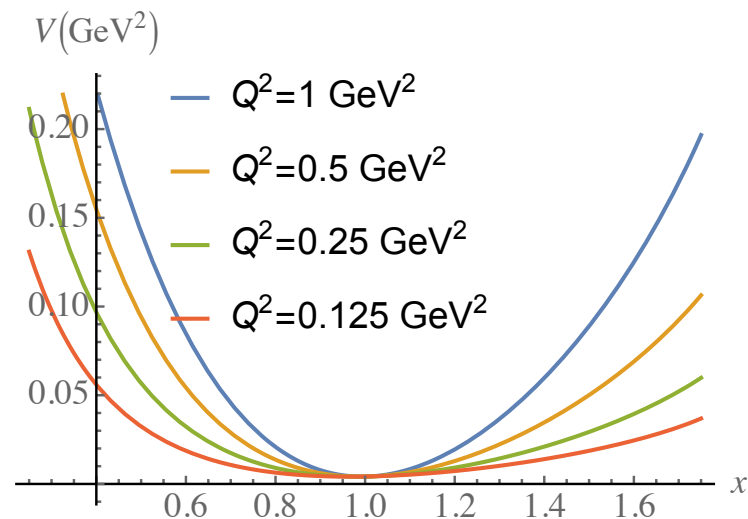
1. Exact
2. On- shell $k^2 = m^2$ in 4-momentum conservation delta function
3. As in 2, Bjorken limit $\frac{Q^2}{\nu^2} \ll 1$

$$\text{Virtuality } V \equiv m^2 - k^2$$



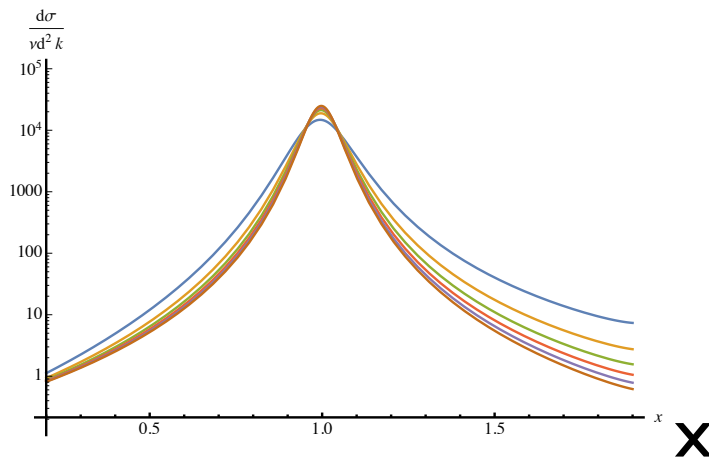
V is very large except for x near 1

Low Q^2 - V is small -
low energy NP

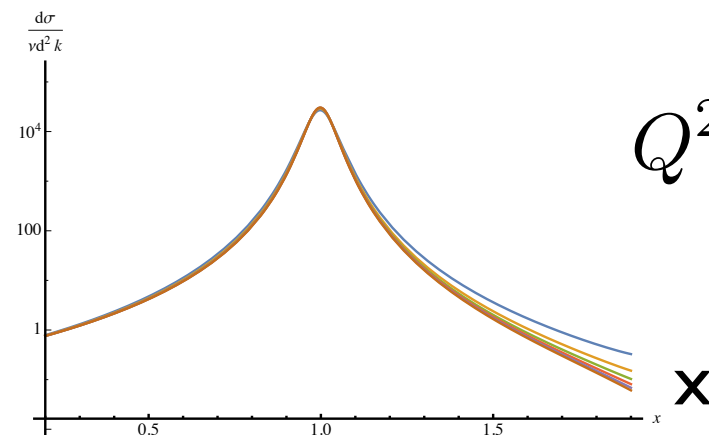


Model exact cross sections

$$\frac{d\sigma}{\nu d^2k}$$

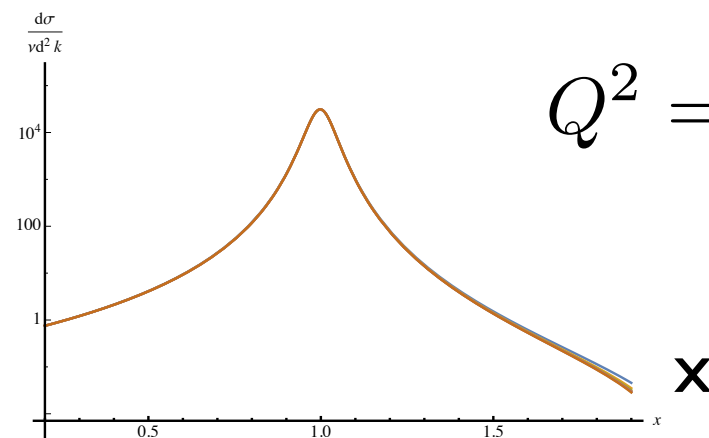


$$Q^2 = 1 - 6 \text{ GeV}^2$$



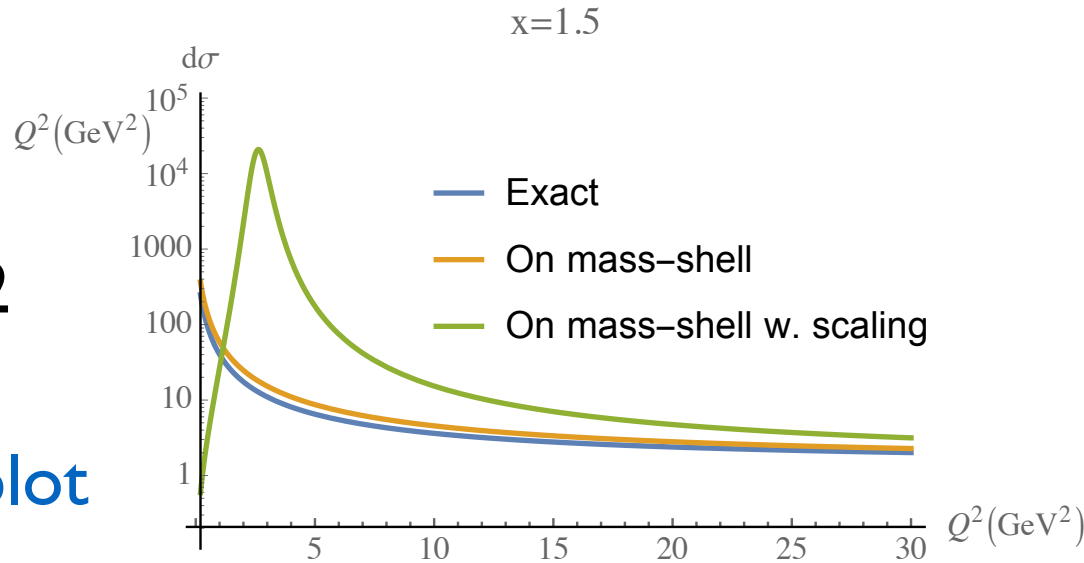
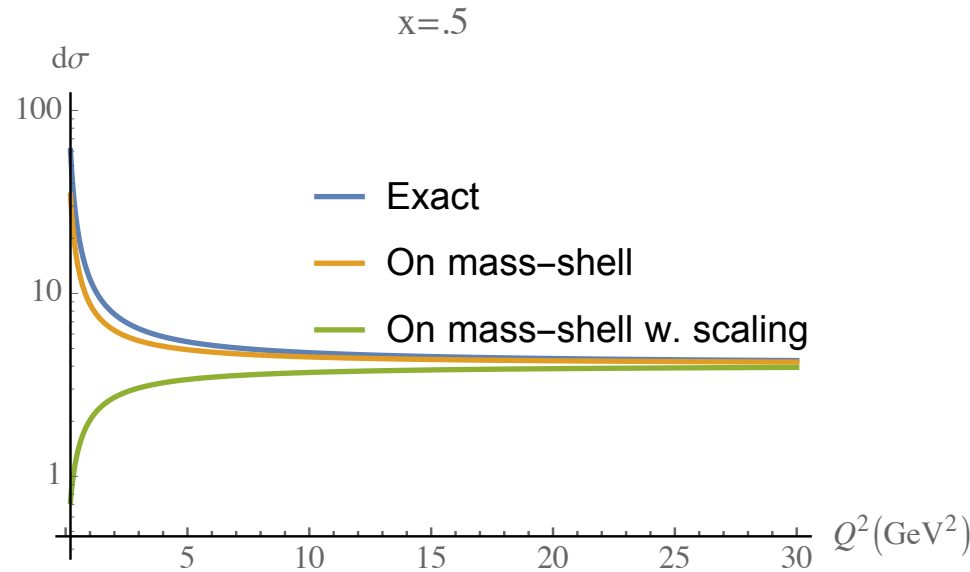
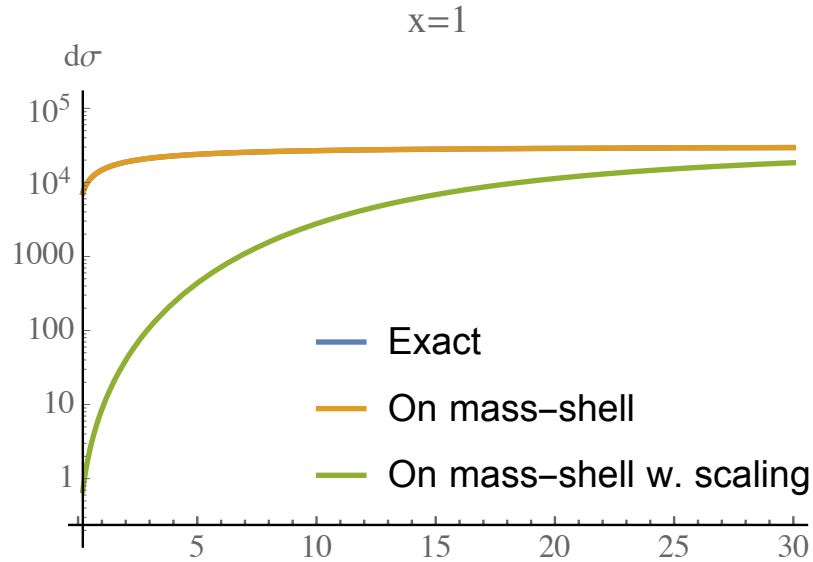
$$Q^2 = 10 - 60 \text{ GeV}^2$$

Scaling at large Q^2



$$Q^2 = 100 - 600 \text{ GeV}^2$$

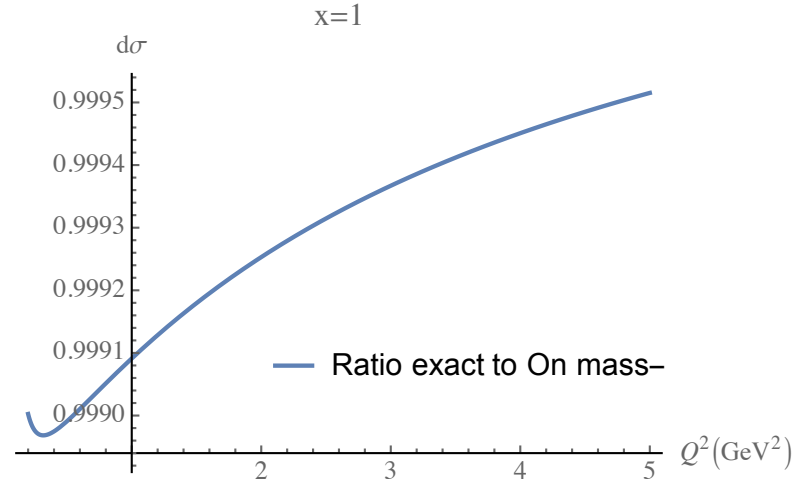
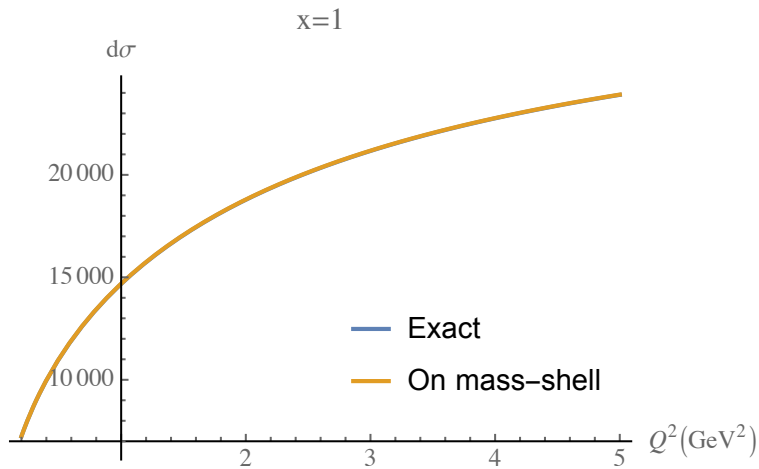
Approach to scaling



Scaling requires high Q^2

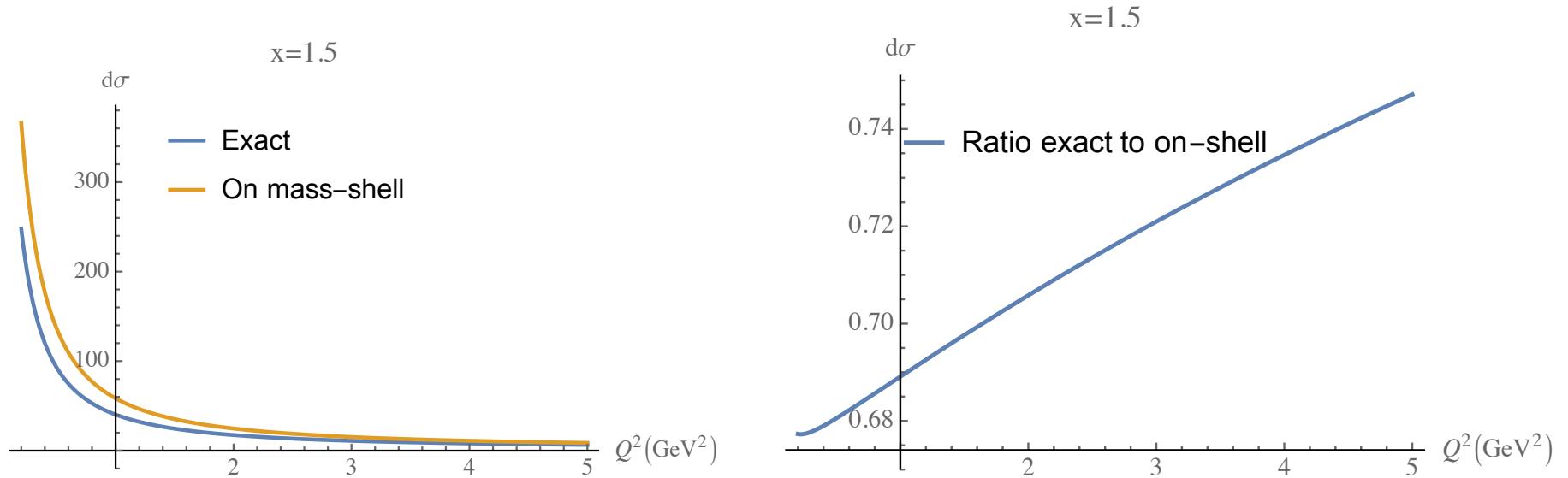
On-shell seems ok in log plot

Low Q , $x=1$ (Long's) exp't



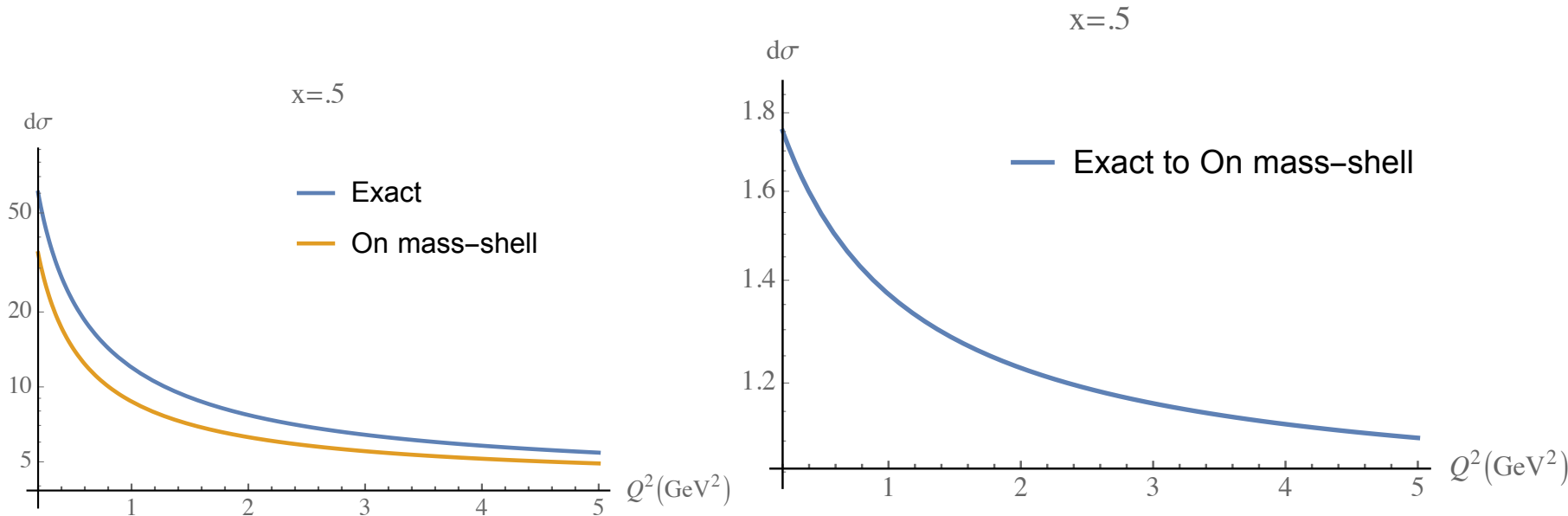
On-mass shell approximation $V=0$ is OK
can us standard wave functions

Low Q, x=1.5



Significant error by assuming on shell-
can't use standard wave functions
spectator on-shell is ok

Low Q, x=0.5

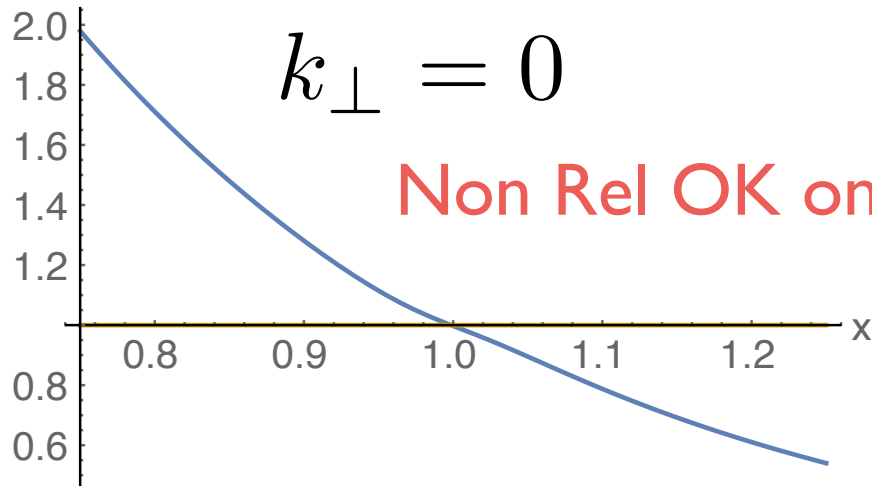


Significant error by assuming on shell
can't use standard wave functions
spectator on-shell is ok

Non-relativistic approximation

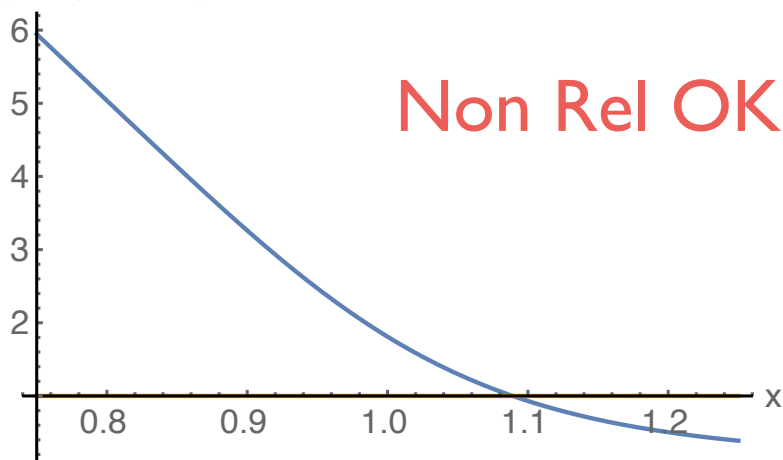
$$Q^2 = 4 \text{ GeV}^2$$

$d\sigma(\text{Exact})/d\sigma(\text{NonRel})$



$$Q^2 = 4 \text{ GeV}^2, k_{\text{perp}} = 0.25 \text{ GeV}/c$$

$d\sigma(\text{Exact})/d\sigma(\text{NonRel})$



Summary

- formulate the NN interaction on the light front - get from BSE
- solve Weinberg equation for the deuteron -done
- prescription from FS 81 review constructs LF wave function from NR wf: study with exact solutions of BSE
- how good is this approx at recoil momenta few hundred MeV?-seems ok up to about 250-300 MeV more study needed
- can we get the LF wf from NN potentials?-seems ok
- Relevance of LF wave functions- OK for space-like form factors
- Quasi-elastic at $Q^2 < 15 \text{ GeV}^2$, LF wave function No Good, Need spectator wf if x is not very near 1
- Non-relativistic approximation is not good except in narrow range of x , depends on k_{perp}

Spares follow

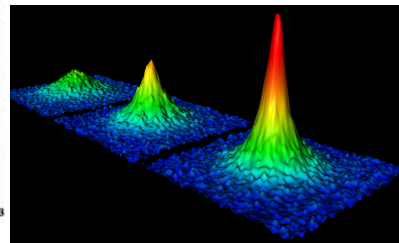
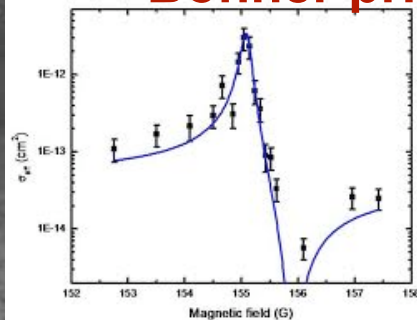
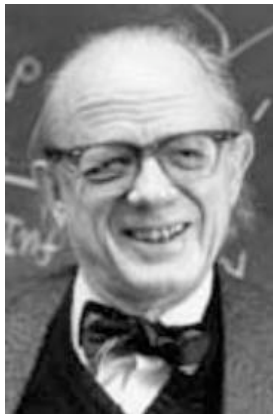
The APS Council and the DNP have endorsed the establishment of the

Herman Feshbach Prize in Nuclear Physics

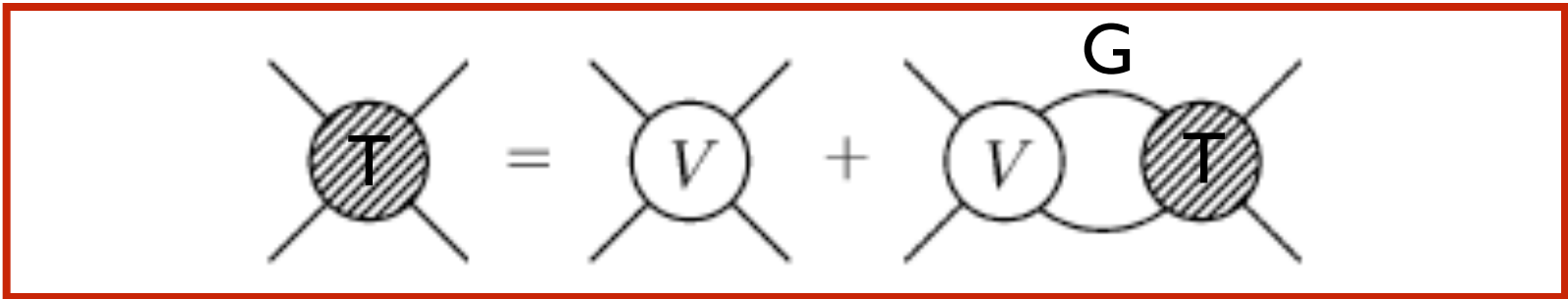
Purpose: To recognize and encourage outstanding research in theoretical nuclear physics. The prize will consist of \$10,000 and a certificate citing the contributions made by the recipient. The prize will be presented biannually or annually.

Herman Feshbach was a dominant force in Nuclear Physics for many years. The establishment of this prize depends entirely on the contributions of institutions, corporations and individuals associated with Nuclear Physics. So far, significant contributions have been made by MIT, the DNP, ORNL/U.Tenn, JSA/SURA, BSA, Elsevier Publishing, TUNL, TRIUMF, MSU, and a number of individuals. More than \$150,000 has been raised, primarily through institutional contributions. **It is very important that physicists make contributions to carry the endowment over the \$200,000 mark, so that the Prize will be eligible to be awarded annually.** Please help us reach that goal by making a contribution. Go online at <http://www.aps.org/> Look for the support banner and click APS member (membership number needed) and look down the list of causes.

If you have any questions, please contact G. A. (Jerry) Miller UW, <mler@uw.edu>.



If annual- number of experimentalists winning Bonner prize goes up by >50%



$$G \sim \frac{1}{\alpha(1-\alpha)} \frac{d^2 p_{\perp} d\alpha}{P^2 - \frac{p_{\perp}^2 + m^2}{\alpha(1-\alpha)}}$$

define p_z : $\frac{p_{\perp}^2 + m^2}{4\alpha(1-\alpha)} = p_{\perp}^2 + p_z^2 + m^2 = \vec{p}^2 + m^2$

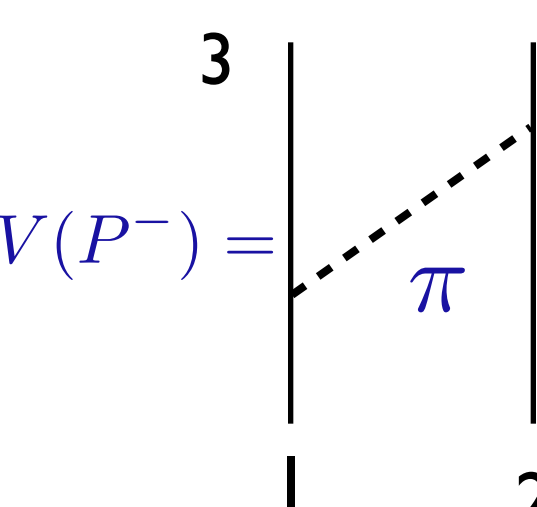
$$G \sim \frac{m^2}{\sqrt{\vec{p}^2 + m^2}} \frac{d^3 p}{p_i^2 - p^2} \text{ Usual propagator with extra factor}$$

Extra factor is close to unity for D wave function

Jason Cooke nucl-th/0112029, Cooke & Miller PRC66, 034002

Solves LF Schroedinger eq (LFSSE)

$$[P_0^- + V(P^-)] |\Psi_D\rangle = P^- |\Psi_D\rangle \quad P^- = 2m - B \quad \text{rest frame}$$



Manifest rotational invar. broken

$$= g^2 \frac{1}{P^- - k_3^- - k_2^- - k_\pi^-}$$

Different meson propagator than Machleidt Miller

2 Solve LSSE using transformation from α to k_z :

$$\alpha = \frac{k^+}{P^+} = \frac{k^+}{2M - B} = \frac{1}{2} \frac{\sqrt{\vec{k}^2 + m^2} + k_z}{\sqrt{\vec{k}^2 + m^2}}$$

Solve w. rot. inv. in \perp plane (polar coords)

Computed B depends on magnetic quantum number!

Cooke nucl-th/0112029, Cooke & Miller PRC66, 034002

Dynamics

- Chiral Lagrangian with $\pi, \eta, \rho, \omega, \delta, \sigma$
- Two meson exchange!
- Explicit P^- dependence

$$K = \text{[Diagram 1]} + \text{[Diagram 2]} + \text{[Diagram 3]} + \text{[Diagram 4]} + \text{[Diagram 5]} + \dots$$

FIG. 1. The first several terms of the full kernel for the Bethe-Salpeter equation of the nuclear model with chiral symmetry.

Cooke nucl-th/0112029, Cooke & Miller PRC66, 034002

Two Meson Dynamics

Instantaneous terms

$$(a) V_{\text{TME:EM}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{1m} - k_{1f} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{2m} - k_{2f} \\ q_i = k_{2m} - k_{2i} \end{array} \right)$$

$$(b) V_{\text{TME:SB}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right)$$

$$(c) V_{\text{TME:SBI}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right) + \left(\begin{array}{c} \text{Diagram 3} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 4} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right)$$

$$(d) V_{\text{TME:SBI}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right)$$

Chiral contact terms

$$(a) V_{\text{TME:C}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2i} - k_{2m} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{1i} - k_{1m} \\ q_i = k_{1i} - k_{1m} \end{array} \right)$$

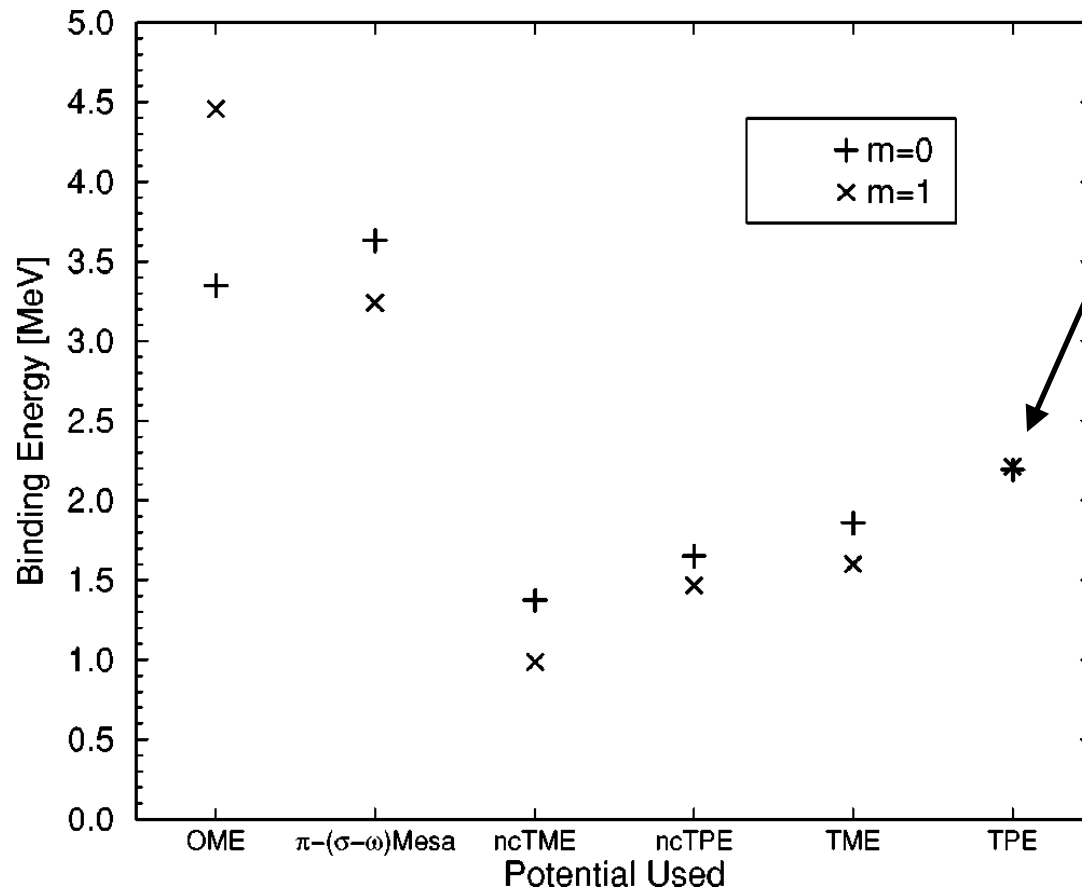
$$(b) V_{\text{TME:SBC}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{2m} - k_{2f} \\ q_i = k_{2i} - k_{2m} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{1m} - k_{1f} \\ q_i = k_{1i} - k_{1m} \end{array} \right) + \left(\begin{array}{c} \text{Diagram 3} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 4} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right)$$

$$(c) V_{\text{TME:SBIC}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f = k_{1f} - k_{1m} \\ q_i = k_{1m} - k_{1i} \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f = k_{2f} - k_{2m} \\ q_i = k_{2m} - k_{2i} \end{array} \right) + \left(\begin{array}{c} \text{Diagram 3} \\ q_f = k_{2m} - k_{2f} \\ q_i = k_{2i} - k_{2m} \end{array} + \begin{array}{c} \text{Diagram 4} \\ q_f = k_{1m} - k_{1f} \\ q_i = k_{1i} - k_{1m} \end{array} \right)$$

$$(d) V_{\text{TME:SBCC}} = \left(\begin{array}{c} \text{Diagram 1} \\ q_f \\ q_i \end{array} + \begin{array}{c} \text{Diagram 2} \\ q_f \\ q_i \end{array} \right)$$

Restoring Rot. Inv.

PRC66, 034002



Uses only
2 pion exch
in 2BE

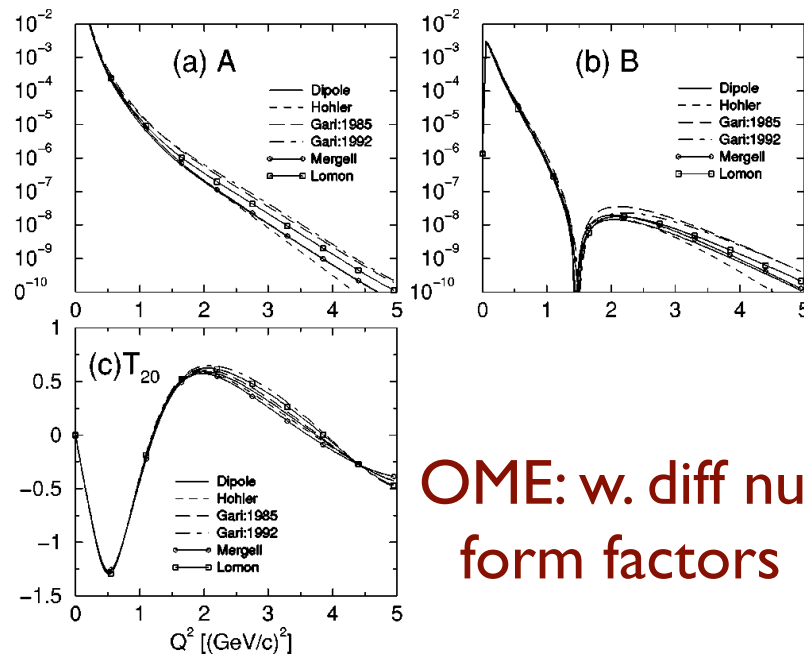
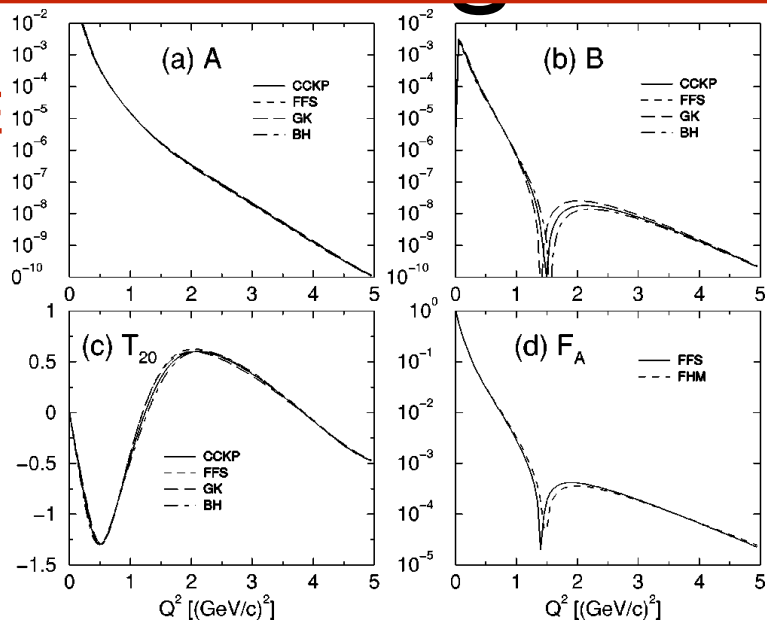
FIG. 9. The values of the binding energy for the $m=0$ and $m=1$ states for different nucleon-nucleon light-front potentials. The σ

Restoring RI in form factors

- Rotational invariance gives angular condition FS
- Angular condition is upheld better when Deut is computed using only one meson exchange OME potentials than two meson exchange TME
- However, form factors do not depend much on choice of bad currents

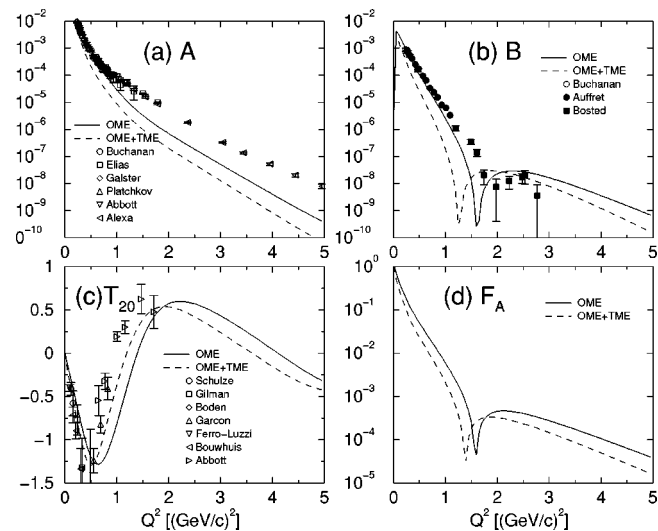
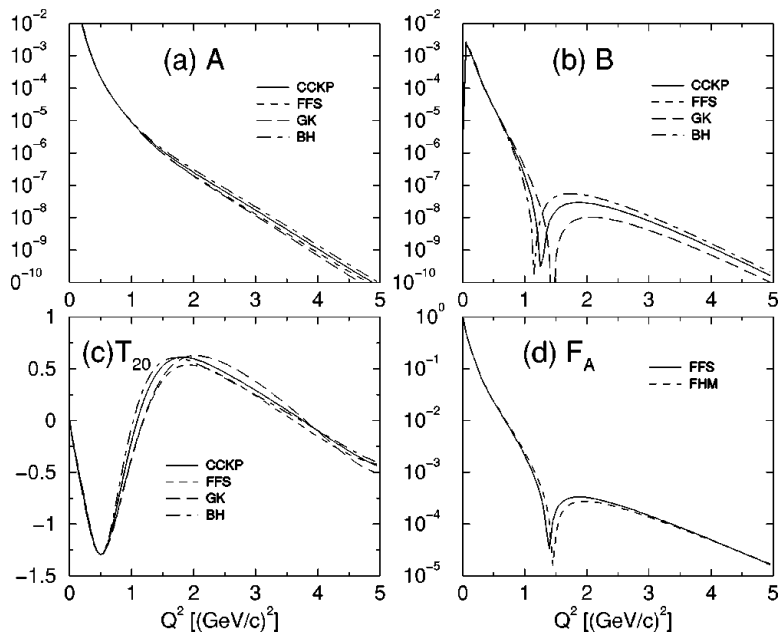
Restoring? RI in form factors

OME



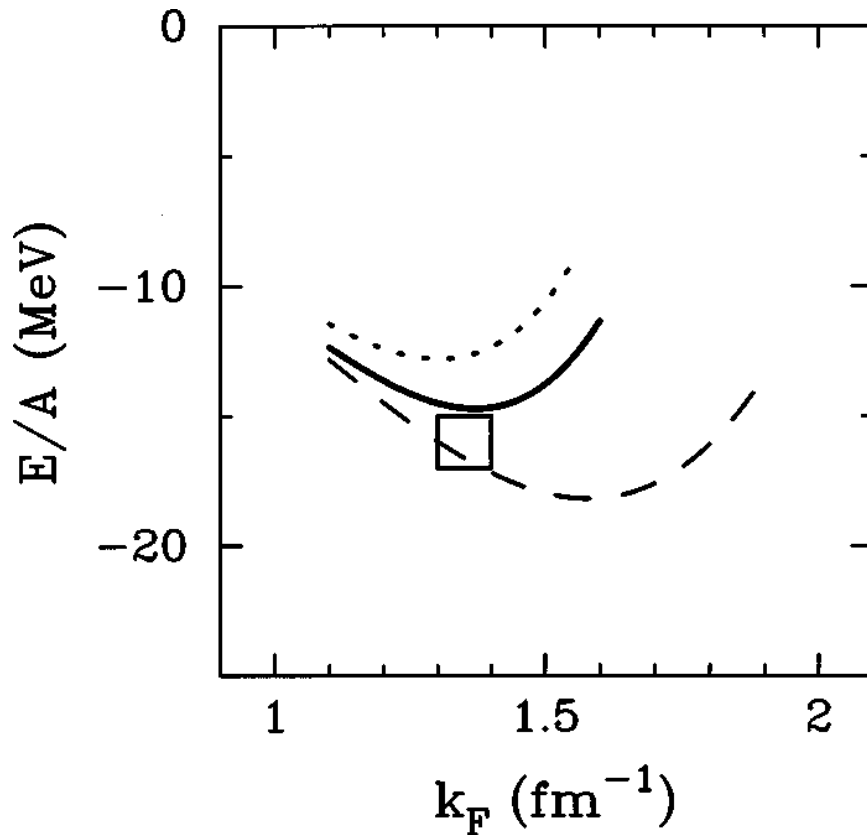
OME: w. diff nucleon form factors

TME



Miller & Machleidt PRC 60,035202

Nuclear Matter Saturation



Solid -our light front
Dashed- ET formalism